# Metric Entropy and Ordinary Differential Equations 

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## 1

Let $C_{1}$ be the set of vector fields $\Phi=\left(\varphi_{1}, \varphi_{2}\right)$ in the plane, each of whose components fulfills the boundedness conditions

$$
\left|\varphi_{i}\right| \leqslant 1, \quad\left|\varphi_{i}\left(z_{1}\right)-\varphi_{i}\left(z_{2}\right)\right| \leqslant\left|z_{i}-z_{2}\right| .
$$

To each vector field $\Phi$ is associated a family of mappings (a flow) $F_{t}$ :

$$
F_{0}(z)=z, \quad \frac{\partial}{\partial t} F_{t}(z)=\Phi\left(F_{t}(z)\right), \quad-\infty<t<\infty .
$$

We denote by $S_{1}$ the set of all mappings $F_{1}$, realized in this way from fields $\Phi$ in $C_{1}$, and by $S_{1}(\delta)$ the set of all mappings obtained under the further condition $\varphi_{1} \geqslant \delta>0$. In the following statement $N_{\epsilon}(S)$ is the number of sets required to cover a set of functions $S$, each set having diameter $<\varepsilon$ in the uniform memtric on the square $Q: 0 \leqslant x \leqslant 1,0 \leqslant y \leqslant 1$.

Theorem. $\quad N_{\epsilon}\left(S_{1}\right) \geqslant \exp c_{1} \varepsilon^{-2}, 0<\varepsilon<\mathrm{I}$,

$$
N_{\epsilon}\left(S_{1}(\delta)\right) \leqslant \exp c_{2}(\delta) \varepsilon^{-4 / 3}, \quad 0<\varepsilon<1 .
$$

The lower bound is very elementary and is included for completeness. Let $g(z)=\left(1-|z|^{2}\right)^{+}$and let $z_{1}, \ldots, z_{N}$ be points in $Q$, with $\left|z_{i}-z_{j}\right| \geqslant \varepsilon$ when $i \neq j$. This can be done so that $N \geqslant c \varepsilon^{-2}$. (From this point onward we use $c$ to denote positive constants.) Now we set $\varphi_{2}=0, \varphi_{1}=\Sigma \pm \varepsilon g\left(\varepsilon^{-1} z-\varepsilon^{-1} z_{i}\right)$. Each center $z_{i}$ moves left or right at least $c \varepsilon$ up to time $y=1$, according to
the choice of $\pm$ : we note that $\varphi_{1}$ never changes sign along any orbit $F_{1}(z)$. This leads to the lower bound $\exp c_{1}$ \&

## 3

In obtaining the upper bound we make several reductions, but these have clearly no effect on the generality. First we suppose that $\varphi_{1}$ and $\varphi_{2}$ are of class $C^{1}\left(R^{2}\right)$ and then we suppose that $\varphi_{1}=1$ and $\varphi_{2}=0$ outside the large square $-2 \leqslant x \leqslant 3,-2 \leqslant y \leqslant 3$. (The second adjustment may require an increase in the Lipschitz constants.) Each mapping $F_{t}$ is differentiable and its Jacobian (or differential) is a matrix $J(t, z)$. The entries of $J(t, z)$ for $0 \leqslant t \leqslant 1$ are uniformly bounded for all $\Phi$ in $C_{1}$. Moreover $J(t, z)$ depends differentiably on $t$, according to the equation of variation (compare $\mid 1$. p. 96|)

$$
(\partial / \partial t) J(t, z)=A\left(F_{t}(z)\right) \cdot J(t, z)
$$

with

$$
A_{i j}(z)=\partial \Phi_{i /} / \partial x_{j} \quad\left(x_{1}=x, x_{2}=y^{\prime}\right)
$$

From the identity

$$
J\left(1, F_{s}(z)\right) \cdot J(s, z)=J(s+1, z)
$$

or

$$
J\left(1, F_{s}(z)\right)=J(s+1, z) J(s, z)^{1}
$$

we see that

$$
\left|\frac{\hat{c}}{\partial t} J\left(1, F_{t}(z)\right)\right| \leqslant c, \quad \text { at } \quad t=0 .
$$

It should be noted that the operator $(\partial / \partial t) f\left(F_{t}(z)\right)$ at $t=0$ equals $L(f)$, and $L=\varphi_{1}(\partial / \partial x)+\varphi_{2}(\partial / \partial y)$. Our inequality can be stated as follows for each component $u$ of $F_{1}(z)$, and each operator $D=\partial / \partial x$ or $\partial / \partial y|L(D u)| \leqslant c$, where $L$ must be defined using the flow $F_{i}$. Now, formally, $L D u-D L u=$ $g_{1}(x, y) u_{x}+g_{2}(x, y) u_{y}$, with $\left|g_{1}\right|+\left|g_{2}\right| \leqslant c$, so hat $L u$ satisfies a uniform Lipschitz condition in the plane. (The proof of the identity for $L D-D L$ proceeds by a smoothing of $\Phi$ so that $u$ is $C^{2}$ and all derivatives have their usual interpretation.)

Let $0<\varepsilon<1$ and let $C_{1}$ be devided into sets $B_{1}$ of diameter $\varepsilon^{2 / 3}$ in the uniform metric of $R^{2}$ (or, what is the same, on the square $-2 \leqslant x \leqslant 3$.
$-2 \leqslant y \leqslant 3$ ). The number of sets necessary here is $\exp c \varepsilon^{-4 / 3}$ (by $\mid 2$, p. 153|). We choose one vector field $\psi_{1}, \psi_{2}$ from some $B_{r}$ and try to estimate the functions $u$ for vector fields $\Phi$ in $B_{r} ;\left|\varphi_{1}-\psi_{1}\right|<\varepsilon^{2 / 3},\left|\varphi_{2}-\psi_{2}\right|<\varepsilon^{2 / 3}$. Henceforth $L=\psi_{1}(\partial / \partial x)+\psi_{2}(\partial / \partial y)$, and $L(\varphi)$ is the operator with $\varphi_{1}, \varphi_{2}$.

We proved before that

$$
\left|\left(L_{\theta} u\right)\left(z_{1}\right)-\left(L_{6} u\right)\left(z_{2}\right)\right| \leqslant c\left|z_{1}-z_{2}\right|,
$$

whence $\left|\left(L_{6} u\right)\left(z_{1}\right)-(L u)\left(z_{2}\right)\right| \leqslant c\left|z_{1}-z_{2}\right|+c \varepsilon^{2 / 3}$. We map the half-plane $t \geqslant 0,-\infty<y<\infty$ to the half-plane $x \geqslant 0$ by the formula $H(t, y)=F_{t}(0, y)$. The rectangle $R: 0 \leqslant t \leqslant \delta^{\prime \prime},|y| \leqslant 1+2 \delta^{-1}$ covers $Q$, and $t, y$ are coordinates of class $C^{\prime}$; the partials of $t, y$ and of the inverse transformation are bounded by a function $\delta$ alone. Henceforth we write $u(t, y)$ in place of $u(H(t, y))$, observing that $\partial / \partial t=L$.

The functions $u(t, y)$ are divided into subsets of diameter $\varepsilon$ in the uniform metric over the set composed of lines $t=m \varepsilon^{1 / 3}$, or rather the segements of these lines contained in $R$; the number of sets needed is $\exp c \varepsilon^{-4 / 3}$ by $\mid 2$, Chap. 10| and elementary properties of the uniform metric.

Suppose now that $m \varepsilon^{1 / 3}<t<(m+1) \varepsilon^{1 / 3}$ and that $t$ has the special form

$$
m \varepsilon^{1 / 3}+\left(\frac{\coprod_{1}^{1}}{1} a_{j} 2^{-j}\right) \varepsilon^{1 / 3}, \quad a_{j}=0 \text { or } 1 .
$$

We write $t_{0}=m \varepsilon^{1 / 3}, t_{j}=m \varepsilon^{1 / 3}+\left(\sum_{1}^{j} a_{j} 2^{-1}\right) \varepsilon^{1 / 3}$ and use finite differences in the formula

$$
u\left(t_{N}, y\right)=u\left(t_{0}, y\right)+\sum_{i}^{N} u\left(t_{j}, y\right)-u\left(t_{j-1}, y\right)
$$

To represent all functions $u\left(t_{N}, y\right)$ we use $<c \varepsilon^{-1 / 3} 2^{j}$ differences $u\left(t_{j}, y\right)-$ $u\left(t_{j 1}, y\right)$. We are going to find the metrical properties of these differences as functions of $y$.

By the mean-value theorem

$$
\begin{aligned}
& u(t+h, y+k)-u(t+h, y)-u(t, y+k)+u(t, y) \\
& \quad=h\left[u_{t}(t+\theta h, y+k)-u_{t}(t+\theta h, y)\right], \quad 0<\theta<1,
\end{aligned}
$$

whence the symmetric difference is at most $c|h||k|+c|h| \varepsilon^{2 / 3}$. We suppose that $h=2^{-j} \varepsilon^{1 / 3}(j=1,2,3, \ldots)$; if $k=2^{j} j^{-2} \varepsilon^{2 / 3}$ the estimate for the symmetric difference is less than $c j^{-2} \varepsilon+c 2^{-j} \varepsilon \leqslant c j^{-2} \varepsilon$. Thus each function $u(t+h, y)-u(t, y)$ varies by at most $c j^{-2} \varepsilon$ when $y$ varies on an interval of length $2^{j} j{ }^{m 2} \varepsilon^{2 / 3}$; clearly $|u(t+h, y)-u(t, y)| \leqslant c|h|$. By the procedure already cited $|2|$, we can cover the set of all such functions by sets of
diameter $c{ }^{2}{ }^{2}$, using no more than $\exp \mathrm{cj}^{2} 2^{{ }^{\prime}} \varepsilon{ }^{2 /} \exp c$ sets. As explained above. this number must be raised to the power $\mathrm{C}^{i} \varepsilon{ }^{1 / 3}$ yielding $\exp c 2^{j} \varepsilon{ }^{13} \exp c j^{2} \varepsilon^{-1}$. This type of estimation is continued until $h=2^{-j} \varepsilon^{1 / 3}<\varepsilon$, or $2^{j}<\varepsilon^{2 / 3}$, so the distance between the lines, on which $u$ is estimated, is at most $\varepsilon$. Thus $2^{j} \leqslant 2 \varepsilon^{23}$. So that multiplication of the estimates for these numbers $j$ yields $<\exp c c^{4 / 3}$ sets. Because $\sum j^{?}$ converges we obtain a covering by sets of diameter $c \varepsilon$.

## 4

The upper bound is somewhat surprising, in as much as $F_{1}$ need not belong to $C^{r}$ for any $r>1$. (When $1<r<2, C^{r}$ is the class of function with derivatives in the Holder-class of exponent $r-1$.) This, however, is not the main obstable to improving the exponent $4 / 3$; if we study fields in which $\varphi_{2}=0, \varphi_{1} \geqslant \delta>0$, we can obtain any exponent $>1$, by the interpolation process used above.

It is natural to look for co-ordinates $u, v$ of class $C^{2}$, in which $\dot{v}=0$, that is $\varphi_{1}(x, y) v_{x}+\varphi_{2}(x, y) v_{y}=0$. This is a first-order partial differential equation. to which we must add $\nabla \ell \neq 0$. In contrast to the behavior of ordinary differential equations, solutions $v$ are generally no smoother than $\Phi$. To illustrate this we take the simplest example: $\varphi_{1} \equiv 1, \varphi_{2}=\varphi_{2}(y)$. We suppose that $\varphi_{2}(y)=0$ for $y \leqslant 0$ and $\varphi_{2} \in C^{\prime}(-\infty, \infty)$, while $\varphi_{2}^{\prime}$ increases on ( 0,1 ). Because $v_{x}=0$ for $y \leqslant 0, v_{y}(0, y)=a \neq 0$. Proceeding formally from the relation $v_{x}+\varphi_{2}(y) v_{y}=0$, we obtain $\left(v_{y}\right)_{x}+\varphi_{2}^{\prime}(y) v_{y}+\varphi_{2}(y) v_{y y}=0$. This has to be interpreted as a derivative along the trajectories of the system $\dot{x}=1, \dot{y}=\varphi_{2}(y)$, as in Section 3, and can be justified if we first treat $\left(v_{y}\right)_{x}$ and $v_{y y}$ as distributions, or generalized derivatives, [3, p. 49]. For economy in writing, we set $\psi=v_{v}$, so $\psi(0, y)=a \neq 0$; we assume as we can, that $a>0$. Let $(x(t), y(t))$ be a solution of the differential system, with initial position $x(0)=0, y(t)=\bar{y}>0$. If $\bar{y}$ is sufficiently small, then $\bar{y} \leqslant y(t) \leqslant$ $2 \bar{y} \leqslant 1$ over $0 \leqslant t \leqslant 1$, because $\varphi_{2}(y)=o(y)$ near $y=0$. Along the solution $\left(x(t) . y^{\prime}(t)\right)$, we have $\dot{\psi}=-\varphi_{2}^{\prime}(y) \psi \leqslant-\phi_{2}^{\prime}(\bar{y}) \psi$, so that $\psi(x(1), y(1)) \leqslant \psi(0, \bar{F})$ $\exp -\varphi_{2}^{\prime}(\bar{r})$. Both the initial point $(0, \bar{y})$ and the point $(1, y(1))$ have distance at most $2 \bar{y}$ from the $x$-axis, on which $\psi=a$. The modulus of continuity of $\psi$. in the square $0 \leqslant x \leqslant 1,0 \leqslant y \leqslant 1$, admits a lower bound $w(\bar{y}) \geqslant c \varphi_{2}^{\prime}(\bar{y})$, and this means that $v_{y}$ is no smoother than $\varphi_{2}^{\prime}$. In particular, when $\varphi_{2}(y)=y^{2}$, for $y \geqslant 0, v$ cannot be of class $C^{2}$.

The inequality $N_{\epsilon}\left(S_{1}(\delta)\right) \geqslant \exp c_{3}(\delta) \varepsilon^{-1}$ is not difficult; we omit the proof because we do not know the best exponent.

# DIFFERENTIAL EQUATIONS AND ENTROPY 

## References

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